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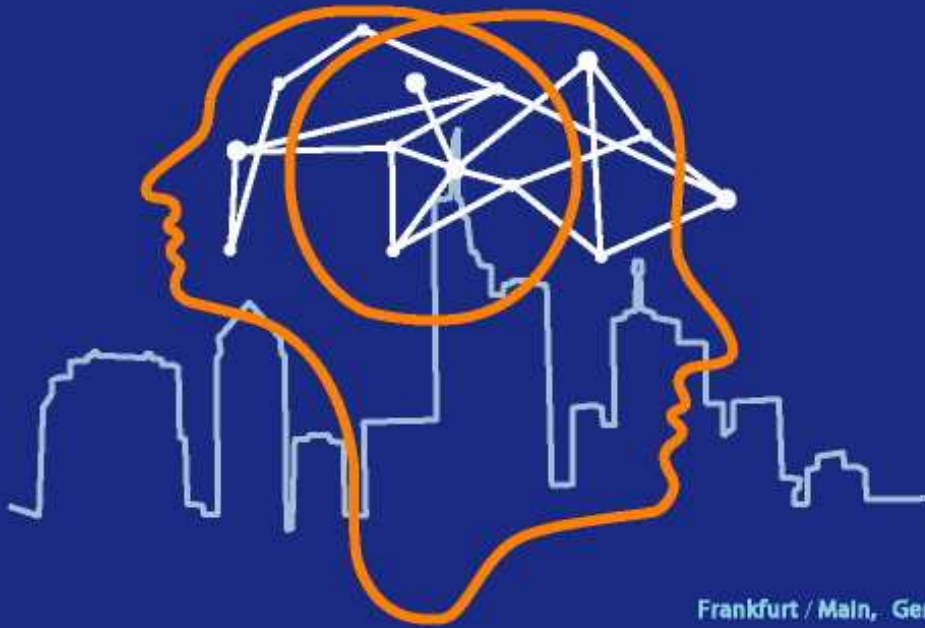
# THEORETICAL

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# NEUROSCIENCE

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# & COMPLEX SYSTEMS



Frankfurt / Main, Germany

**Pre-School  
Linear Algebra**

Cornelius Weber



## Determinant

The idea is to characterize a square matrix by a real number. So one defines a function  $\det : M(n \times n) \rightarrow \mathbb{R}$  from an  $n \times n$ -matrix to a real number. The following three axioms lead to a unique definition of such a function:

### D1

$$\det(\mathbb{1}) = 1$$

( $\det$  is normalized.)

### D2

if  $A$  has two similar rows, then  $\det(A) = 0$

( $\det$  is alternating.)

### D3

a) if for three row vectors  $a, b, c$  holds  $a + b = c$ , then:

$$\det \begin{pmatrix} \vdots \\ c \\ \vdots \end{pmatrix} = \det \begin{pmatrix} \vdots \\ a \\ \vdots \end{pmatrix} + \det \begin{pmatrix} \vdots \\ b \\ \vdots \end{pmatrix}$$

b) if for two row vectors  $a, b$  holds  $a = \lambda b$ , then:

$$\det \begin{pmatrix} \vdots \\ a \\ \vdots \end{pmatrix} = \lambda \det \begin{pmatrix} \vdots \\ b \\ \vdots \end{pmatrix}$$

( $\det$  is linear in every row.)

The following follows from these axioms:

- if the matrix has a row  $a$  where  $a = (0, \dots, 0)$ , then  $\det(A) = 0$ .  
This follows from axiom 3.
- if we can construct  $B$  from  $A$  by adding the  $j$ th row to the  $i$ th row,  $j \neq i$ , then  $\det(B) = \det(A)$ .  
To see this, partition  $B$  according to axiom 3 into two matrices, then one will have two similar rows and hence  $\det = 0$  according to axiom 2.
- if  $A$  is diagonal, then  $\det A = a_{11} \cdot a_{22} \cdot \dots \cdot a_{nn}$ .  
To see this, start with the identity matrix (axiom 1) and apply operations of axiom 3.
- $\det(A) = 0 \Leftrightarrow$  the row vectors (also column vectors) are linearly dependent.  
(Has to do with axiom 2: via elementary row operations, bring the matrix into block diagonal form. Then  $\det(A)$  is the product over the pivot elements. It holds: vectors linearly independent  $\Leftrightarrow$  pivot elements  $\neq 0$ .)

Furthermore:

$$\bullet \det \begin{pmatrix} a \\ \vdots \\ b \end{pmatrix} = -\det \begin{pmatrix} b \\ \vdots \\ a \end{pmatrix} \quad (\det \text{ is anti-symmetric})$$

- $\det(AB) = \det(A) \cdot \det(B)$
- but:  $\det(A + B) \neq \det(A) + \det(B)$ , in general.
- $A$  invertible  $\Leftrightarrow \det(A) \neq 0$

## Application to a homogeneous system of equations

$$A\vec{x} = 0$$

has a non-trivial solution ( $\vec{x} \neq 0$ ) only if  $\det(A) = 0$ .

Reason: consider  $\det(A) \neq 0 \Leftrightarrow$  (column) vectors are linearly independent. Hence, there is *no* linear combination of column vectors (that's what is done by multiplying  $A$  from the right with a vector  $\vec{x}$ ), which fulfills  $A\vec{x} = 0$  (except the trivial solution  $\vec{x} = 0$ ).

So in order to allow for a non-trivial solution of  $\vec{x}$ , it has to be  $\det(A) = 0$ . The solution  $\vec{x}$  will not be unique, because  $c\vec{x}$  will also be a solution.

## Application to an inhomogeneous system of equations

$$A\vec{x} = \vec{b}$$

has a *unique* solution iff  $\det(A) \neq 0$ .

Reason:  $\det(A) \neq 0 \Leftrightarrow$  vectors are linearly independent, and form basis vectors of the vector space. Each point  $\vec{b}$  can be reached by unique coefficients  $\vec{x}$  of these basis vectors.

## Geometrical interpretation

The determinant of a matrix is the area of the parallelogram spanned by its row (column-) vectors (neglecting the determinant's sign).

To see this, one transforms the parallelogram via two kinds of operations into a cube. The first kind makes a rectangular box out of the parallelogram while keeping the volume constant. This is done by adding  $\lambda$  times the  $j$ th row vector  $\vec{v}_j$  to the  $i$ th row vector  $\vec{v}_i$ ,  $j \neq i$ . That does not change the determinant (see above), and also not the volume (see Figure).

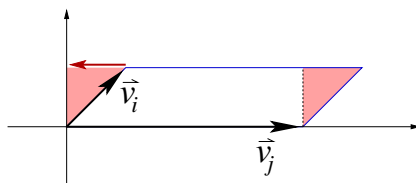


Figure 1: Varying  $\vec{v}_i$  by a portion of  $\vec{v}_j$  shifts the shaded volume of the parallelogram.

Now our parallelogram has right angles. The second kind of operations scales each vector to unity length. The determinant – just as the volume – scales linear with the scaling of each of these vectors. These scaling factors tell us about the original volume.

Resulting is a cube with volume 1. We might have to rotate it to align with the axis, which neither changes the volume, nor the determinant since a rotation matrix has determinant 1.

Another interpretation: if one applies the matrix to the points of a cube, then the determinant gives the distortion of the volume. A determinant which is smaller than 0 means that the volume has changed its orientation (the order of the axes permutes by an odd number).

## Calculation of the determinant (Leibnitz' rule)

$$\det \begin{pmatrix} a_1 \\ \vdots \\ a_n \end{pmatrix} \stackrel{\mathbf{D3}}{=} \sum_{i_1=1}^n a_{1i_1} \cdot \det \begin{pmatrix} e_{i_1} \\ a_2 \\ \vdots \\ a_n \end{pmatrix} \stackrel{\mathbf{D1}}{=} \sum_{i_1=1}^n \sum_{i_2=1}^n \cdots \sum_{i_n=1}^n a_{1i_1} \cdot a_{2i_2} \cdot \dots \cdot a_{ni_n} \cdot \det \begin{pmatrix} e_{i_1} \\ \vdots \\ e_{i_n} \end{pmatrix}$$

The first step uses **D3a)** by constructing a row from unity vectors, as well as **D3b)** by multiplying with the corresponding value.

### Example of Leibnitz' rule

$$\begin{aligned} \det \begin{pmatrix} a & b \\ c & d \end{pmatrix} &= a \det \begin{pmatrix} 1 & 0 \\ c & d \end{pmatrix} + b \det \begin{pmatrix} 0 & 1 \\ c & d \end{pmatrix} \\ &= a c \det \begin{pmatrix} 1 & 0 \\ 1 & 0 \end{pmatrix} + a d \det \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + b c \det \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + b d \det \begin{pmatrix} 0 & 1 \\ 0 & 1 \end{pmatrix} \\ &= a d - b c \end{aligned}$$

Source: Gerd Fischer Lineare Algebra P. 142 (in german language)

## The Eigenvalue Equation

Let  $V$  be a  $M$ -dimensional vector space  $\mathbb{R}$ . Let  $A: V \rightarrow V$  be a linear function ( $A$  can be expressed as a  $M \times M$ -matrix). A real number  $\lambda_m$  is an eigenvalue of  $A$ , if there exists a vector  $\vec{v}_m \in V$  with  $\vec{v}_m \neq 0$ , so the eigenvalue equation holds:

$$A\vec{v}_m = \lambda_m\vec{v}_m \tag{1}$$

Hence, the  $\vec{v}_m$  are those vectors, which are extended/compressed, but **not** rotated, by the function  $A$ . If  $\vec{v}_m$  is an eigenvector, then  $c\vec{v}_m$  is also an eigenvector. The eigenvalue equation for  $M$  eigenvalues and eigenvectors in matrix notation is:

$$AU = U\Lambda \tag{2}$$

or

$$\begin{pmatrix} & & & \\ & A & & \\ & & & \end{pmatrix} \begin{pmatrix} \left( \begin{matrix} \vec{v}_1 \\ \cdot \\ \cdot \end{matrix} \right) \cdots \left( \begin{matrix} \vec{v}_M \\ \cdot \\ \cdot \end{matrix} \right) \end{pmatrix} = \begin{pmatrix} \left( \begin{matrix} \vec{v}_1 \\ \cdot \\ \cdot \end{matrix} \right) \cdots \left( \begin{matrix} \vec{v}_M \\ \cdot \\ \cdot \end{matrix} \right) \end{pmatrix} \begin{pmatrix} \lambda_1 & & 0 \\ & \ddots & \\ 0 & & \lambda_M \end{pmatrix}$$

$\Lambda$  is a diagonal matrix, with the eigenvalues on the diagonal.  $U$  is a matrix, with the eigenvectors as its columns. Mind the order on both sides of the equation!

In Eq. 2 we assume that there are  $M$  eigenvectors. This sounds reasonable, because in an  $M$ -dimensional vector space there are maximally  $M$  linear independent vectors.

In the following, we will address the following questions:

- How to compute eigenvalues & eigenvectors?
- What is an inverse matrix?
- How to invert a matrix?

Let  $U^{-1}U = \mathbb{1}_M$  (here  $U^{-1}$  is the inverse matrix of  $U$ ) then we can write:

$$U^{-1}AU = \Lambda \quad \Leftrightarrow \quad AU = U\Lambda \quad \Leftrightarrow \quad A = U\Lambda U^{-1}.$$

Left: Eq. 2 left-multiplied by  $U^{-1}$

Right: Eq. 2 right-multiplied by  $U^{-1}$ .

## Relevance

We have:  $\vec{y} = A\vec{x}$  ← the elements of  $\vec{x}$  are combined in a complicated matter;  
an element of  $\vec{y}$  is made from many elements of  $\vec{x}$

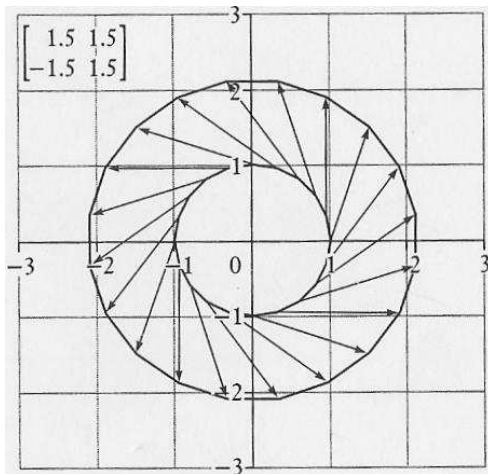
We want:  $\tilde{y} = \Lambda\tilde{x}$  ← each element of  $\tilde{x}$  is multiplied by an eigenvalue;  
an element of  $\tilde{y}$  is made from only one element of  $\tilde{x}$

How to get there:  $\vec{y} = U\Lambda U^{-1}\vec{x}$

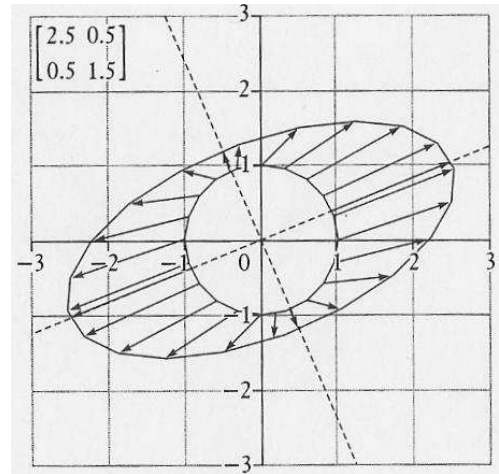
Left-multiply by  $U^{-1}$ :  $U^{-1}\vec{y} = \Lambda U^{-1}\vec{x}$

Define:  $\tilde{y} := U^{-1}\vec{y}$  and  $\tilde{x} := U^{-1}\vec{x}$

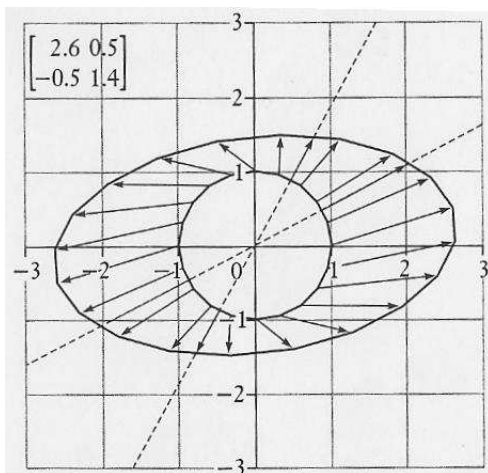
Back-transform:  $\vec{y} = U\tilde{y}$  and  $\vec{x} = U\tilde{x}$



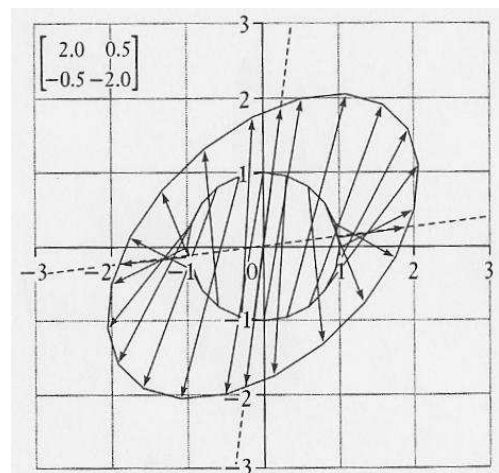
rotation matrix  
no eigenvectors



symmetric matrix  
orthogonal eigenvectors



two positive eigenvalues



one positive, one negative eigenvalue

Figure 1: Effect of  $2 \times 2$ -Matrices on sample vectors.

Source: "Consider the Lowly 2x2 Matrix" by Alexander Kujath: <http://www2.in.tu-clausthal.de/~hormann/teaching/ProSemSS05/PS.CG.24.05.2005.a.ppt>

**Orthonormal:**  $M$  vectors  $\in \mathbb{R}^M$  are orthonormal if

$$\vec{v}^i \cdot \vec{v}^j = \delta_{ij}$$

In matrix notation:

$$\begin{pmatrix} (\dots \vec{v}^1 \dots) \\ \vdots \\ (\dots \vec{v}^M \dots) \end{pmatrix} \left( \begin{pmatrix} \vec{v}^1 \\ \vdots \\ \vdots \end{pmatrix} \cdots \begin{pmatrix} \vec{v}^2 \\ \vdots \\ \vdots \end{pmatrix} \right) = \begin{pmatrix} 1 & & 0 \\ & \ddots & \\ 0 & & 1 \end{pmatrix}$$

or

$$SS^T = \mathbb{1}_M \tag{3}$$

which is equivalent to (by left-multiplying with  $S^{-1}$ )

$$S^T = S^{-1}$$

The column vectors of  $S^T$  form an orthonormal base of  $\mathbb{R}^M$ . Hence, the eigenvectors would span an orthonormal base, if  $U$  (on the previous page) is orthonormal. Note that orthogonal eigenvectors may always be scaled to be orthonormal.

Source: Gerd Fischer Lineare Algebra P. 200 (in german language)

**When are the eigenvectors orthogonal?** We will show that of any symmetric matrix, eigenvectors belonging to different eigenvalues are orthogonal.

Let  $A\vec{v}_j = \lambda_j\vec{v}_j$  and  $A\vec{v}_i = \lambda_i\vec{v}_i$  (i.e.  $\vec{v}_j$  is eigenvector belonging to eigenvalue  $\lambda_j$  and  $\vec{v}_i$  is eigenvector belonging to eigenvalue  $\lambda_i$ ). Then:

$$\lambda_j\vec{v}_i^T\vec{v}_j = \vec{v}_i^T\lambda_j\vec{v}_j = \vec{v}_i^T A\vec{v}_j = (A^T\vec{v}_i)^T\vec{v}_j \stackrel{\text{symm}}{=} (A\vec{v}_i)^T\vec{v}_j = \lambda_i\vec{v}_i^T\vec{v}_j$$

If  $\lambda_j \neq \lambda_i$  then it has to be  $\vec{v}_i^T\vec{v}_j = 0$ .

Source: Gerd Fischer Lineare Algebra P. 208 (in german language)

**How find eigenvalues & eigenvectors?** The eigenvalue equation (1) is equivalent to:

$$(A - \lambda_m\mathbb{1})\vec{v}_m = 0 \tag{4}$$

This system of equations for  $\vec{v}_m$  has non-trivial solutions if

$$\det(A - \lambda_m\mathbb{1}) = 0 \tag{5}$$

Computing this determinant results in a polynomial in  $\lambda_m$ , the *characteristic polynomial* of  $A$ . Its solutions are the eigenvalues  $\{\lambda_m\}$  of  $A$ .

The eigenvectors  $\vec{v}_m$  are then found by inserting the found  $\lambda_m$  into equation (4). Since the rows of this matrix are linearly dependent ( $\det = 0$ ), the solutions are not unique. One obtains only a relation which fixes the direction but not the length of the eigenvectors.

**When are the eigenvalues real?** As an example, let us look only at the characteristic polynomial of a  $2 \times 2$ -matrix.

$$\begin{aligned} \det \begin{pmatrix} a_{11} - \lambda & a_{12} \\ a_{21} & a_{22} - \lambda \end{pmatrix} \\ = \lambda^2 - a_{11}\lambda - a_{22}\lambda + a_{11}a_{22} - a_{21}a_{12} = \lambda^2 + (-a_{11} - a_{22})\lambda + a_{11}a_{22} - a_{21}a_{12} \end{aligned}$$

which we set to zero to obtain the eigenvalues. We apply the quadratic formula

$$x_{1/2} = -\frac{b}{2} \pm \sqrt{\frac{b^2}{4} - c}$$

where we identify  $b = (-a_{11} - a_{22})$  and  $c = a_{11}a_{22} - a_{21}a_{12}$ . The eigenvalues are *both* real if:  $c \leq \frac{b^2}{4}$ , i.e.

$$a_{11}a_{22} - a_{21}a_{12} \leq \frac{(-a_{11} - a_{22})^2}{4} = \frac{a_{11}^2 + a_{22}^2 + 2a_{11}a_{22}}{4}$$

$$0 \leq \frac{a_{11}^2 + a_{22}^2 + 2a_{11}a_{22} - 4a_{11}a_{22} + 4a_{21}a_{12}}{4}$$

$$0 \leq a_{11}^2 + a_{22}^2 - 2a_{11}a_{22} + 4a_{21}a_{12}$$

$$0 \leq (a_{11} - a_{22})^2 + 4a_{21}a_{12}$$

This is in particular true if the off-diagonal elements have the same sign.

In general one can show: A symmetric matrix has real eigenvalues.

**When are all eigenvalues positive?** If the matrix is [positive definite](#) – later.

**A matrix that has no eigenvectors** A rotation matrix in  $\mathbb{R}^2$  is given by

$$A = \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix}$$

The characteristic polynomial  $P_A$  is

$$\lambda^2 - 2 \cos \alpha \cdot \lambda + \cos^2 \alpha + \sin^2 \alpha = \lambda^2 - 2 \cos \alpha \cdot \lambda + 1$$

We set this to zero and apply the quadratic formula to get the eigenvalues:

$$x_{1/2} = -\frac{b}{2} \pm \sqrt{\frac{b^2}{4} - c} \longrightarrow \lambda_{1/2} = \cos \alpha \pm \sqrt{\cos^2 \alpha - 1} \stackrel{\sin^2 + \cos^2 = 1}{=} \cos \alpha \pm \sqrt{-\sin^2 \alpha}$$

The term under the square root is non-negative only for  $\alpha = 0, \pi$ , etc. Then the matrix is similar to the identity matrix.

(Defining the complex number  $i = \sqrt{-1}$  we have  $\lambda_{1/2} = \cos \alpha \pm i \sin \alpha$ .)

Source: Gerd Fischer Lineare Algebra P. 169 (in german language)

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## Derivation of the quadratic formula

$$\begin{aligned} x^2 + bx + c &= 0 \\ x^2 + bx &= -c \\ x^2 + bx + \frac{b^2}{4} &= \frac{b^2}{4} - c \quad \text{this step is motivated by the following} \\ \left(x + \frac{b}{2}\right)^2 &= \frac{b^2}{4} - c \\ x_{1,2} + \frac{b}{2} &= \pm \sqrt{\frac{b^2}{4} - c} \\ x_{1,2} &= -\frac{b}{2} \pm \sqrt{\frac{b^2}{4} - c} \end{aligned}$$

## What is an inverse matrix?

An  $M \times M$  matrix  $A$  is invertible, if there exists an *inverse matrix*  $A^{-1}$  with

$$A A^{-1} = \mathbb{1}_M$$

Since  $A^{-1} = A^{-1} \mathbb{1}_M = A^{-1}(A A^{-1}) = (A^{-1} A) A^{-1}$  it will also be

$$A^{-1} A = \mathbb{1}_M$$

## How to invert a matrix?

**1. Method:** For a given matrix  $S$  search the inverse unknown matrix  $X = S^{-1}$ . Ansatz:

$$S X = \mathbb{1}$$

For each of the  $M$  column vectors of  $X$  this is a system of equations with  $M$  unknowns (the components of the column vector) and  $M$  equations (the rows of  $S$ ). The solution of each of these systems of equations results in a column of  $X$ .

Source: Bronstein Taschenbuch der Mathematik P. 206 (in german language)

**2. Method:** With **elementary matrices**  $B$  one can transform every invertible matrix  $S$  into the identity matrix  $\mathbb{1}$ . In the following, apply these to  $S$  (left side) so that it becomes  $\mathbb{1}$ , and apply the same operations to  $\mathbb{1}$  (right side) so that it becomes the inverse  $S^{-1}$  of  $S$ :

$$B_n \cdot \dots \cdot B_1 \cdot S = \mathbb{1} \quad \Leftrightarrow \quad B_n \cdot \dots \cdot B_1 \cdot \mathbb{1} = S^{-1}$$

Source: Gerd Fischer Lineare Algebra P. 94 (in german language)

**What are elementary matrices?** Every invertible matrix  $S$  can be expressed as a product of elementary matrices. These are square matrices which differ from  $\mathbb{1}$ -matrix in only few entries. They perform – from left multiplied – elementary row operations:

- (i) multiplying a row by  $\lambda$ ,
- (ii) adding  $\lambda$  times the  $j$ -th row to the  $i$ -th row and
- (iii) interchanging two rows.

$A^{-1}$  is symmetric if  $A$  is symmetric, written also as  $A^{-1} = (A^{-1})^T$ . Reason:

$$\mathbb{1}_M = A^{-1} A = (A^{-1} A)^T = A^T (A^{-1})^T = \underbrace{A (A^{-1})^T}_{\text{must be } A^{-1}}$$

Source: Christopher M. Bishop Neural Networks for Pattern Recognition Appendix A

$A$  and  $A^{-1}$  have the same eigenvectors; the eigenvalues are reciprocal. Reason:

$$\begin{array}{l|l} A \vec{v} = \lambda \vec{v} & \text{left-multiply with } A^{-1} \\ \vec{v} = \lambda A^{-1} \vec{v} & \text{divide by } \lambda \\ \lambda^{-1} \vec{v} = A^{-1} \vec{v} & \leftarrow \text{eigenvalue equation for the inverse matrix} \end{array}$$

Source: Christopher M. Bishop Neural Networks for Pattern Recognition Appendix A

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## Taylor Series

Assume one can write a function  $f : \mathbb{R} \rightarrow \mathbb{R}$  as a polynomial:

$$f(x) = a_0 + a_1x + a_2x^2 + a_3x^3 + \dots + a_nx^n$$

The Taylor formula states, if this polynomial series is valid in an interval, the following: The function value at any point  $x$  within this interval can be computed if one knows all derivatives of a function at (another) point  $x_0$  which is also in this interval.

$$f(x) = a_0 + a_1(x - x_0) + a_2(x - x_0)^2 + a_3(x - x_0)^3 + \dots + a_n(x - x_0)^n$$

(For  $x_0 = 0$  this yields the topmost equation.)

Now we need the coefficients  $\{a_k\}$ . We can obtain these by comparing the coefficients above to the values obtained from deriving the function at  $x_0$ . The derivatives at  $x_0$  are:

$$\begin{aligned} f'(x) &= a_1 + 2a_2(x - x_0) + 3a_3(x - x_0)^2 + \dots + na_n(x - x_0)^{n-1} \\ f^{(2)}(x) &= 2a_2 + 6a_3(x - x_0) + \dots + n(n-1)a_n(x - x_0)^{n-2} \\ f^{(3)}(x) &= 6a_3 + \dots + n(n-1)(n-2)a_n(x - x_0)^{n-3} \end{aligned}$$

At our expansion point  $x = x_0$  we have:

$$\begin{aligned} f(x_0) = a_0 &\Leftrightarrow a_0 = f(x_0) \\ f'(x_0) = a_1 &\Leftrightarrow a_1 = f'(x_0) \\ f^{(2)}(x_0) = 2a_2 &\Leftrightarrow a_2 = \frac{f^{(2)}(x_0)}{2!} \\ f^{(3)}(x_0) = 6a_3 &\Leftrightarrow a_3 = \frac{f^{(3)}(x_0)}{3!} \\ f^{(n)}(x_0) = n!a_n &\Leftrightarrow a_n = \frac{f^{(n)}(x_0)}{n!} \end{aligned}$$

With these coefficients  $a_k$  the Taylor series can be written as:

$$f(x) = \sum_{k=0}^{\infty} a_k(x - x_0)^k = \sum_{k=0}^{\infty} \frac{f^{(k)}(x_0)}{k!}(x - x_0)^k$$

Examples (for the last one we need  $i^2 = -1$ ):

$$\begin{aligned} e^x &= 1 + x + \frac{1}{2!}x^2 + \frac{1}{3!}x^3 + \frac{1}{4!}x^4 + \frac{1}{5!}x^5 + \dots = \sum_k^{\infty} \frac{x^k}{k!} \\ \cos x &= 1 - \frac{1}{2!}x^2 + \frac{1}{4!}x^4 - \dots = \sum_k^{\infty} \frac{x^{2k}}{(2k)!}(-1)^k \\ \sin x &= x - \frac{1}{3!}x^3 + \frac{1}{5!}x^5 - \dots = \sum_k^{\infty} \frac{x^{2k+1}}{(2k+1)!}(-1)^k \\ e^{ix} &= 1 + ix - \frac{1}{2!}x^2 - i\frac{1}{3!}x^3 + \frac{1}{4!}x^4 + i\frac{1}{5!}x^5 - \dots = \sum_k^{\infty} \frac{ix^k}{k!} \\ &= \cos x + i \sin x \end{aligned}$$

The latter is the *Euler Formula*.

From the Euler Formula follows:

- (i)  $e^{ix} + e^{-ix} = \cos x + i \sin x + \cos x - i \sin x = 2 \cos x$
- (ii)  $e^{ix} - e^{-ix} = \cos x + i \sin x - (\cos x - i \sin x) = 2i \sin x$
- (iii)  $\cos^2 x + \sin^2 x = \frac{e^{2ix} + e^{-2ix} + 2 - e^{2ix} - e^{-2ix} + 2}{4} = 1$
- (iv)  $\cos(\varphi + \psi) + i \sin(\varphi + \psi) = e^{i(\varphi + \psi)} = e^{i\varphi} e^{i\psi} = (\cos \varphi + i \sin \varphi)(\cos \psi + i \sin \psi) = \cos \varphi \cos \psi - \sin \varphi \sin \psi + i(\sin \varphi \cos \psi + \cos \varphi \sin \psi)$

From these series one can furthermore see:

$$\cosh x := \frac{e^x + e^{-x}}{2} = \sum_k^{\infty} \frac{x^{2k}}{(2k)!} = \cos ix \quad \text{or} \quad \cos x = \cosh ix$$

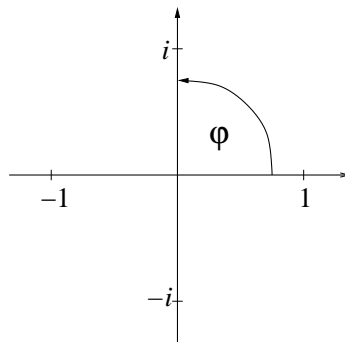
$$\sinh x := \frac{e^x - e^{-x}}{2} = \sum_k^{\infty} \frac{x^{2k+1}}{(2k+1)!} = -i \sin ix \quad \text{or} \quad \sin x = -i \sinh ix$$

$$\cosh^2 x - \sinh^2 x = \frac{e^{2x} + e^{-2x} + 2}{4} - \frac{e^{2x} + e^{-2x} - 2}{4} = 1$$

## Complex Numbers

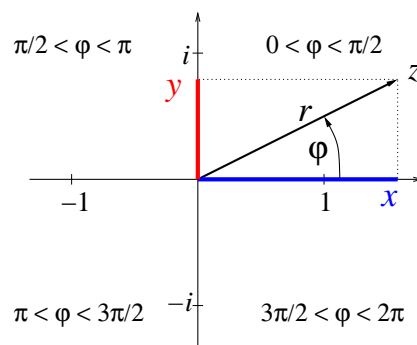
Problem: we cannot find a solution to a negative square root. The solution is to define:

$$i := \sqrt{-1}$$



Squaring  $i$  means doubling its angle. So multiplications are somehow defined by rotations. Then e.g.  $(-i)^2 = -1$ , but  $-i \cdot i = 1$ .

Real numbers are a subset of complex numbers, where the angle is a multiple of  $\pi$ . Complex numbers are isomorph to  $\mathbb{R}^2$ .



$$\begin{aligned} z &= x + iy \\ &= r \cos \varphi + ir \sin \varphi = r e^{i\varphi} \end{aligned}$$

The latter is the Euler formula (verified by comparing Taylor series of  $\cos \varphi$ ,  $\sin \varphi$  and  $e^{i\varphi}$ ). Addition of complex numbers ...

$$\begin{aligned} z_1 + z_2 &= \overbrace{x_1 + x_2 + i(y_1 + y_2)} \\ &= r_1 e^{i\varphi_1} + r_2 e^{i\varphi_2} \end{aligned}$$

... is vector addition (which we can better see with the Cartesian coordinates).

Multiplication of complex numbers ...

$$\begin{aligned} z_1 \cdot z_2 &= x_1 x_2 + i x_1 y_2 + i x_2 y_1 - 1 \cdot y_1 y_2 = x_1 x_2 - y_1 y_2 + 2i x_1 y_2 \\ &= r_1 e^{i\varphi_1} \cdot r_2 e^{i\varphi_2} = \underbrace{r_1 r_2}_{\text{absolute values}} e^{i(\varphi_1 + \varphi_2)} \end{aligned}$$

... multiplies the absolute values and adds the angles (better seen in the polar coordinates).

## Revisiting a matrix without (real) eigenvectors

$$A = \begin{pmatrix} a & b \\ -b & a \end{pmatrix}$$

The characteristic polynomial  $P_A$  is

$$\det \begin{pmatrix} a - \lambda & b \\ -b & a - \lambda \end{pmatrix} = \lambda^2 - 2a\lambda + a^2 + b^2 \stackrel{!}{=} 0$$

Applying the quadratic formula

$$\lambda_{1/2} = a \pm \sqrt{-b^2} = \underbrace{a \pm ib} = r(\cos \varphi + i \sin \varphi) = r e^{\pm i\varphi}$$

Indeed, in the original matrix we could have already expressed  $(a, b)$  as  $(r \cos \varphi, r \sin \varphi)$  for some suitable  $(r, \varphi)$ . Hence, we have a “rotation matrix” that also scales the vector by  $r$ . We insert the given eigenvalues (in the notation of the under-braced term) into the eigenvalue equation to yield the eigenvectors. For  $\lambda_1 = a + ib$ :

$$\begin{aligned} av_1 - av_1 - ibv_1 + bv_2 &= 0 &\Rightarrow v_2 &= iv_1 \\ -bv_1 + av_2 - av_2 - ibv_2 &= 0 \end{aligned}$$

For  $\lambda_2 = a - ib$ :

$$\begin{aligned} av_1 - av_1 + ibv_1 + bv_2 &= 0 &\Rightarrow v_2 &= -iv_1 \\ -bv_1 + av_2 - av_2 + ibv_2 &= 0 \end{aligned}$$

The two found eigenvectors are

$$\begin{pmatrix} 1 \\ i \end{pmatrix}, \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

Note that each of these vectors has 4 entries (2 of which are zero here), so vectors are in a 4-dimensional space. Any complex multiples of these vectors are also eigenvectors.

Interpretation of the eigenvalue equation in complex space:

- each component of the eigenvector is scaled by the absolute value of the eigenvalue
- the direction of each component is turned by the angle argument of the eigenvalue

Hence, the direction of the eigenvectors does not remain constant in complex space.

## Potential Method

Finding eigenvalues (and eigenvectors) can be difficult since there is no formula for solving a (characteristic) polynomial of higher than 3rd order. A simple way to find the eigenvector belonging to the largest eigenvalue of a symmetric matrix  $A$  is the following iterative procedure.

First, choose an initial vector  $\vec{v}$ . Then repeat:

1. apply  $A\vec{v}$
2. normalize  $\vec{v}$  (e.g.  $l^2$  or  $l^\infty$  norm)

until converged. The eigenvalue is then trivially obtained.

Source: <http://alf.math.uni-rostock.de/kurt/wiwi/slides/wiwi-15.pdf>

## Correlation Based Learning

Let us consider a linear neuron which computes its activation  $y$  from an input  $\vec{x}$  as

$$y = \sum_j w_j x_j$$

The Hebbian Rule for the change  $\Delta w_{ki}$  of a weight (synaptic efficacy) between a neuron  $i$  with activity  $x_i$  of an input layer and a neuron  $k$  with activity  $y_k$  of the output layer is

$$\Delta w_{ki} = \epsilon \cdot x_i \cdot y_k = \epsilon \cdot x_i \cdot \sum_j w_{kj} x_j$$

where  $\epsilon$  is a small learning step size. Let us express the input activations  $\{x_i\}$  as a vector  $\vec{x}$ :

$$\Delta \vec{w}_k = \epsilon \vec{x} y_k = \epsilon \vec{x} (\vec{w}_k^\top \cdot \vec{x}) = \epsilon \vec{x} \vec{x}_k^\top \vec{w} \stackrel{\text{averaged}}{=} \epsilon \langle \vec{x} \vec{x}_k^\top \rangle \vec{w} =: \epsilon C \vec{w}$$

$C$  is the correlation matrix of the input patterns (stimuli) in the input layer.

## First order derivatives in higher dimensions

### f: $\mathbb{R} \mapsto \mathbb{R}$

The derivative of a function  $f: \mathbb{R} \mapsto \mathbb{R}$  is the slope of the tangent. This tangent is expressed in the Taylor series

$$f(x_0 + \xi) = f(x_0) + a\xi + \dots$$

by the first term (after the constant), namely the linear function  $a \cdot \xi$ . There,  $a = \frac{df(x)}{dx}$  is the first order derivative of  $f$ .

### f: $\mathbb{R}^N \mapsto \mathbb{R}$

Let  $f: \mathbb{R}^N \mapsto \mathbb{R}$ , then the derivative is given by a tangent (hyper-) plane: along each of the  $N$  coordinate axes of the input space there may be a different slope. We can write these values into a vector, the *gradient*  $\vec{a}$  of  $f$ . The variation of the function value around the point of expansion  $\vec{x}_0$  depends on the direction of the vector of variation  $\vec{\xi}$ . This is expressed by the dot product  $\vec{a} \cdot \vec{\xi} = \vec{a}^T \vec{\xi} = \sum_{i=1}^N a_i \xi_i$ :

$$f(\vec{x}_0 + \vec{\xi}) = f(\vec{x}_0) + \vec{a}^T \vec{\xi} + \dots \quad (i = 1, \dots, N)$$

### f: $\mathbb{R}^N \mapsto \mathbb{R}^M$

Let  $f: \mathbb{R}^N \mapsto \mathbb{R}^M$ , so for each of the  $M$  components of the function value (a vector) we have such an equation:

$$f_j(\vec{x}_0 + \vec{\xi}) = f_j(\vec{x}_0) + \vec{a}_j^T \vec{\xi} + \dots \quad (j = 1, \dots, M)$$

We write these  $M$  equations below each other to form one matrix equation

$$\vec{f}(\vec{x}_0 + \vec{\xi}) = \vec{f}(\vec{x}_0) + A\vec{\xi} + \dots$$

So the first order term of a Taylor series in a higher dimensional vector space is represented by a matrix, since a matrix represents a linear operator.

The  $M \times N$ -matrix  $A$  is called *Differential* or *Jacobi-Matrix*. An entry  $a_{ji}$  is the partial derivative of the  $j$ -th component of the function value w.r.t. the  $i$ -th component of the input vector:

$$a_{ji} = \frac{\partial f_j}{\partial x_i}(\vec{x}_0) = \lim_{(\vec{\xi})_i \rightarrow 0} \frac{f_j(\vec{x}_0 + (\vec{\xi})_i) - f_j(\vec{x}_0)}{(\vec{\xi})_i}$$

## 2nd order derivative

**f:  $\mathbb{R}^N \mapsto \mathbb{R}^1$**

Let  $f: \mathbb{R}^N \mapsto \mathbb{R}^1$ , then we can express the Taylor series up to 2nd order easily:

$$\begin{aligned} f(\vec{x}_0 + \vec{\xi}) &= f(\vec{x}_0) + \sum_{i=1}^N \frac{\partial f}{\partial x_i} \xi_i + \frac{1}{2} \sum_i^N \sum_j^N \frac{\partial^2 f}{\partial x_i \partial x_j} \xi_i \xi_j + \dots \\ &= f(\vec{x}_0) + \vec{\nabla} f \cdot \vec{\xi} + \frac{1}{2} \vec{\xi}^T H \vec{\xi} + \dots \end{aligned}$$

The first order derivative yields the gradient vector  $\vec{\nabla} f$ . The second order derivative is represented by the *Hesse matrix*. The *Hesse Matrix* is symmetric, because  $\frac{\partial^2 f}{\partial x_i \partial x_j} = \frac{\partial^2 f}{\partial x_j \partial x_i}$ .

**Explanation for f:  $\mathbb{R}^2 \mapsto \mathbb{R}^1$**

Let  $f: \mathbb{R}^2 \mapsto \mathbb{R}^1$ , then the change of a function value along the direction of a vector  $\vec{\xi}$  is composed of its change along each of the vector's components,  $\xi_1$  and  $\xi_2$ :

$$f(x_1^0 + \xi_1, x_2^0 + \xi_2) = f(x_1^0, x_2^0) + \underbrace{\frac{\partial f}{\partial x_1} \Big|_{(x_1^0, x_2^0)}}_{g(x_1^0, x_2^0)} \xi_1 + \underbrace{\frac{\partial f}{\partial x_2} \Big|_{(x_1^0, x_2^0)}}_{h(x_1^0, x_2^0)} \xi_2 \quad (1)$$

To get 2nd-order precision we must consider that the derivatives (here defined as  $g$  and  $h$ ) themselves change along  $\vec{\xi}$ :

$$\begin{aligned} \Delta g &:= g(\vec{x}^0 + \vec{\xi}) - g(\vec{x}^0) = \frac{\partial g}{\partial x_1} \xi_1 + \frac{\partial g}{\partial x_2} \xi_2 \\ \Delta h &:= h(\vec{x}^0 + \vec{\xi}) - h(\vec{x}^0) = \frac{\partial h}{\partial x_1} \xi_1 + \frac{\partial h}{\partial x_2} \xi_2 \end{aligned}$$

This expresses how much the derivatives have changed along the whole extent of  $\vec{\xi}$ . The contributions to the change of the function value  $f$  are only the mean of these changes:  $g + \frac{1}{2} \Delta g$  and  $h + \frac{1}{2} \Delta h$ .

So in Eq. 1 we have to replace  $g \rightarrow g + \frac{1}{2} \Delta g$  and  $h \rightarrow h + \frac{1}{2} \Delta h$ . We obtain the additional terms:

$$\frac{1}{2} \Delta g \xi_1 + \frac{1}{2} \Delta h \xi_2 = \frac{1}{2} \left( \left( \frac{\partial^2 f}{\partial x_1^2} \xi_1 + \frac{\partial^2 f}{\partial x_2 \partial x_1} \xi_2 \right) \xi_1 + \left( \frac{\partial^2 f}{\partial x_1 \partial x_2} \xi_1 + \frac{\partial^2 f}{\partial x_2^2} \xi_2 \right) \xi_2 \right)$$

which is in matrix notation (using  $\frac{\partial^2 f}{\partial x_2 \partial x_1} = \frac{\partial^2 f}{\partial x_1 \partial x_2}$ ):

$$\frac{1}{2} (\xi_1 \ \xi_2) \begin{pmatrix} \frac{\partial^2 f}{\partial x_1^2} & \frac{\partial^2 f}{\partial x_1 \partial x_2} \\ \frac{\partial^2 f}{\partial x_1 \partial x_2} & \frac{\partial^2 f}{\partial x_2^2} \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} =: \frac{1}{2} \vec{\xi}^T H \vec{\xi}$$

with the Hesse matrix  $H$ .

## Positive definite

Definition: A symmetric matrix  $A$  is positive definite if  $\vec{x}^T A \vec{x} > 0 \quad \forall \vec{x}$  with  $\vec{x} \neq 0$ .

Theorem:  $A$  positive definite  $\Leftrightarrow$  all eigenvalues  $> 0$ .

Proof:

“ $\Rightarrow$ ” for the eigenvectors we have:

$$\vec{v}^T A \vec{v} = \vec{v}^T \lambda \vec{v} = \lambda \underbrace{\vec{v}^T \vec{v}}_{>0} \stackrel{!}{>} 0 \Rightarrow \lambda > 0$$

“ $\Leftarrow$ ” express the vector  $\vec{x}$  as a linear combination of eigenvectors:  $\vec{x} = S^{-1} \vec{\alpha}$ . ( $S^{-1}$  is the matrix, where the column vectors are the eigenvectors of  $A$ . These are orthogonal – orthonormal if we wish – in case of a symmetric matrix, if the eigenvalues are different; hence such a linear combination can always be made.) So we can write

$$\vec{x}^T A \vec{x} = (S^{-1} \vec{\alpha})^T A S^{-1} \vec{\alpha} = \vec{\alpha}^T S^{-1T} A S^{-1} \vec{\alpha} = \vec{\alpha}^T \underbrace{S^{-1T} S^{-1}}_{\mathbb{1}} \Lambda \vec{\alpha} = \sum_i \lambda_i \alpha_i^2 > 0$$

The under-braced term assumes orthonormal eigenvectors.  $\Lambda$  is the diagonal matrix of the eigenvalues.

## When are matrices positive definite?

- If  $\vec{y} = A \vec{x}$  then  $\vec{x}^T A \vec{x} = \vec{x}^T \vec{y}$  is a dot product of  $\vec{x}$  with  $\vec{y}$ . This being positive means that the angle between both vectors is less than  $90^\circ$ . Thus, a positive definite matrix does not turn a vector by more than  $90^\circ$ .
- $A, B$  positive definite  $\Rightarrow$  the matrix  $A + B$  also. Reason:

$$\vec{x}^T (A + B) \vec{x} = \underbrace{\vec{x}^T A \vec{x}}_{>0} + \underbrace{\vec{x}^T B \vec{x}}_{>0}$$

- $A$  positive definite  $\Rightarrow A^r$  positive definite, where  $r \in \mathbb{R}$ .  
Reason: eigenvalues  $\lambda_i$  are positive  $\Rightarrow \lambda_i^r$  positive.
- Let  $R$  be any matrix with linearly independent columns (not necessarily square matrix). Then  $R^T R$  is positive definite.  
Reason:  $\vec{x}^T R^T R \vec{x} = (R \vec{x})^T (R \vec{x}) = \|R \vec{x}\|^2$   
This is always positive. The term  $R \vec{x}$  which is a linear combination of the columns of  $R$  cannot become zero for  $\vec{x} \neq 0$ , because the columns are linearly independent.
- $R$  as above, further let  $A$  be positive definite.  $\Leftrightarrow R^T A R$  is positive definite.

## Relevance

Let  $A$  be a symmetric  $2 \times 2$ -matrix.

$$f(\vec{x}) := \begin{pmatrix} x_1 & x_2 \end{pmatrix} \begin{pmatrix} a_{11} & a_{12} \\ a_{21} & a_{22} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = a_{11}x_1^2 + a_{22}x_2^2 + 2a_{12}x_1x_2$$

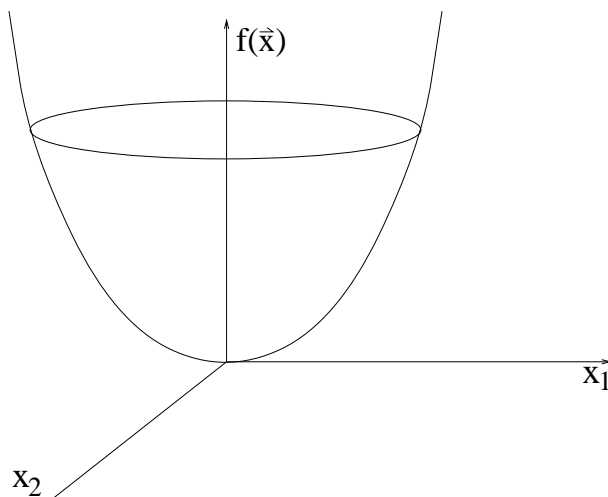
This might be the 2nd order term of a Taylor series. We have

- $f(0) = 0$
- $\frac{\partial f}{\partial x_1} = 2a_{11}x_1 + 2a_{12}x_2$
- $\frac{\partial f}{\partial x_2} = 2a_{22}x_2 + 2a_{12}x_1$

So the first derivative at the origin is zero:  $\left. \frac{\partial f}{\partial x_{1,2}} \right|_{\vec{x}=0} = 0$

- $\frac{\partial^2 f}{\partial x_1^2} = 2a_{11}$ ,  $\frac{\partial^2 f}{\partial x_2^2} = 2a_{22}$ ,  $\frac{\partial^2 f}{\partial x_1 \partial x_2} = \frac{\partial^2 f}{\partial x_2 \partial x_1} = 2a_{12}$

If  $A$  is positive definite, then  $f(\vec{x}) > 0$  for all  $\vec{x} \neq 0$ . So  $f(\vec{x})$  looks like a parabola centred at the origin.



Source: Gilbert Strang Introduction to Applied Mathematics P. 17-19

**Def:**  $(\text{Hesse } f)(\vec{x})_{ij} = \frac{\partial^2 f}{\partial x_i \partial x_j}$ , Hesse matrix, symmetric (“2nd derivative matrix”).

**Theorem:** Let  $f : V \mapsto \mathbb{R}$  be 2x continuously differentiable and  $\vec{\nabla} f(x_o) = 0$ . Then:

If  $(\text{Hesse } f)(\vec{x})$  positive (**negative**) definite, then  $f$  has in  $x_o$  an isolated minimum (**maximum**).

Source: Otto Forster Analysis 2 P. 61/62